q = 0 Oscillator Algebra as a Hopf Algebra

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Received May 25, 1994

In this paper we show that the $q = 0$ oscillator algebra is a Hopf algebra.

About 4 years ago Greenberg (1990) presented an example of infinite statistics for identical particles. He averaged the boson algebra

$$
aa^+ - a^+a = 1 \tag{1}
$$

and fermion algebra

$$
aa^+ + a^+a = 1 \tag{2}
$$

to get the new algebra

$$
aa^+ = 1 \tag{3}
$$

He showed that the statistical mechanics of particles satisfying the algebra (3) obeys the quantum Boltzmann statistics. After this work was done he found out that the algebras (1) - (3) are special examples of a q-deformed boson algebra (Biedenharn, 1989; Macfarlane, 1989; Arik and Coon, 1976),²

$$
aa^{+} - qa^{+}a = 1
$$

[N, a^{+}] = a^{+}
[N, a] = -a (4)

where N is a number operator.

0020-7748/95/0300-0301\$07.50/0 @ 1995 Plenum Publishing Corporation

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²When $q = 1, -1, 0$, the algebra (4) reduces to the algebras (1), (2), and (3), respectively.

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Greenberg assumed the existence of a unique vacuum state annihilated by the annihilation operator a :

$$
a|0\rangle = 0 \tag{5}
$$

He constructed the number operator N for the algebra (3) satisfying the second and third relations of (4), whose form is given by

$$
N = a^{+}a + a^{+}a^{+}aa + a^{+}a^{+}a^{+}aaa + \cdots
$$

=
$$
\sum_{m=1}^{\infty} (a^{+})^{m} a^{m}
$$
 (6)

Here we check that the number operator (6) satisfies the second and third relations of (4):

$$
[N, a^{+}] = \sum_{m=1}^{\infty} (a^{+})^{m} a^{m} a^{+} - \sum_{m=1}^{\infty} (a^{+})^{m+1} a^{m}
$$

\n
$$
= \sum_{m=1}^{\infty} (a^{+})^{m} a^{m-1} a a^{+} - a^{+} \sum_{m=1}^{\infty} (a^{+})^{m} a^{m}
$$

\n
$$
= a^{+} \left(\sum_{m=1}^{\infty} (a^{+})^{m-1} a^{m-1} - \sum_{m=1}^{\infty} (a^{+})^{m} a^{m} \right)
$$

\n
$$
= a^{+} \left(1 + \sum_{m=1}^{\infty} (a^{+})^{m} a^{m} - \sum_{m=1}^{\infty} (a^{+})^{m} a^{m} \right)
$$

\n
$$
= a^{+} \left(1 + \sum_{m=1}^{\infty} (a^{+})^{m} a^{m} - \sum_{m=1}^{\infty} (a^{+})^{m} a^{m} \right)
$$

\n
$$
= a^{+} \qquad (7)
$$

which implies that the number operator given in (6) fulfills the properties for the number operator. This kind of number operator cannot be obtained by setting $q = 0$ in the number operator given in the q-boson algebra. The number operator for the q -boson algebra (4) is already known from

$$
[N] = \frac{q^N - 1}{q - q^{-1}} = a^+ a \tag{8}
$$

Substituting $q = 0$ in equation (7) leads to no relation between the number operator N and mode operators a and a^+ , because a q-number goes to 1 whenever q goes to 0.

Now we will show that this algebra ($q = 0$ oscillator algebra) is a Hopf algebra. The three operations of the Hopf algebra for the $q = 0$ oscillator algebra are given by

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$$
\Delta(N) = N \otimes I + I \otimes N
$$

\n
$$
\Delta(I) = I \otimes I
$$

\n
$$
\Delta(a^+) = I \otimes a^+
$$

\n
$$
\Delta(a) = I \otimes a
$$

\n
$$
\epsilon(I) = 1, \qquad \epsilon(a) = a, \qquad \epsilon(a^+) = a^+, \qquad \epsilon(N) = 0
$$

\n
$$
S(I) = 1, \qquad S(a) = a, \qquad S(a^+) = a^+, \qquad S(N) = -N
$$

where Δ , S, and ϵ are called the coproduct, antipode, and counit, respectively. The Hopf algebra A should satisfy the following three axioms for all elements (generators) $x \in A$:

$$
(\mathrm{id} \otimes \Delta)\Delta(x) = (\Delta \otimes \mathrm{id})\Delta(x)
$$

\n
$$
m(\mathrm{id} \otimes S)\Delta(x) = m(S \otimes \mathrm{id})\Delta(x) = \epsilon(x) \cdot 1
$$

\n
$$
(\epsilon \otimes \mathrm{id})\Delta(x) = (\mathrm{id} \otimes \epsilon)\Delta(x) = x, \qquad x = I, a, a^*, N
$$

It is easy to check that this set of definitions fulfills the three axioms. First let us prove that the coproduct axiom holds for a :

$$
(\mathrm{id}\otimes\Delta)\Delta(a)=(\Delta\otimes\mathrm{id})\Delta(a)
$$

Proof. LHS = (id $\otimes \Delta \Delta(a) = (\mathrm{id} \otimes \Delta)(I \otimes a) = I \otimes \Delta(a) = I \otimes I$ $I \otimes a$; RHS = $(\Delta \otimes id)\Delta(a) = (\Delta \otimes I)(I \otimes a) = \Delta(I) \otimes a = I \otimes I \otimes a$.

Second let us check the case for N:

$$
(\mathrm{id} \otimes \Delta)\Delta(N) = (\Delta \otimes \mathrm{id})\Delta(N)
$$

Proof. LHS = (id $\otimes \Delta$)($N \otimes I + I \otimes N$) = $N \otimes I \otimes I + I \otimes N \otimes I$ $I + I \otimes I \otimes N$; RHS = $(\Delta \otimes id)(N \otimes I + I \otimes N) = N \otimes I \otimes I + I \otimes I$ $N\otimes I + I \otimes I \otimes N$.

The counit axiom is verified as follows:

 $(\epsilon \otimes id)\Delta(a) = (id \otimes \epsilon)\Delta(a) = a$

Proof. LHS = ($\epsilon \otimes id$) $\Delta(a) = (\epsilon \otimes id)(I \otimes a) = \epsilon(I) \otimes a = a$; RHS = $(\Delta \otimes \epsilon)\Delta(a) = (\mathrm{id} \otimes \epsilon)(I \otimes a) = I \otimes \epsilon(a) = a$.

We will prove that

$$
m(\mathrm{id})\Delta(a) = m(S \otimes \mathrm{id})\Delta(a) = \epsilon(a)
$$

Proof. LHS = $m(id \otimes S)\Delta(a) = I \cdot S(a) = a = \epsilon(a)$; RHS = $m(S \otimes$ id) $\Delta(a) = S(I)a = a = \epsilon(a)$.

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We now verify that

 $\Delta(aa^+) = \Delta(a)\Delta(a^+)$

Proof. LHS = $\Delta(aa^+) = \Delta(I) = I \otimes I$; RHS = $(I \otimes a)(I \otimes a^+) = I \otimes I$ $aa^+ = I \otimes I$.

It is easy to check that the following formula holds:

$$
\Delta([N, a]) = [\Delta(N), \Delta(a)]
$$

Proof. $[\Delta(N), \Delta(a)] = [N \otimes I + I \otimes N, I \otimes a] = N \otimes a - N \otimes a +$ $I \otimes Na - I \otimes aN = I \otimes [N, a] = -I \otimes a = -\Delta(a) = \Delta(-a) = \Delta([N, a]).$

From the above proofs, we can conclude that $q = 0$ oscillator algebra is really a Hopf algebra.

ACKNOWLEDGMENTS

This paper was supported in part by the Non Directed Research Fund, Korea Research Foundation (1994); the present studies were supported in part by the Basic Science Research Program, Ministry of Education, 1994 (BSRI-94-2413).

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